

Persistent Currents in Normal Metal Rings: Old Questions, New Answers

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Abstract:

The Aharonov-Bohm effect, first described in 1959, is among the most spectacular effects of quantum mechanics, emphasizing the role the electromagnetic potentials – and not the electromagnetic fields – play for the wave-like motion of quantum particles. Considering a ringlike geometry in a constant perpendicular magnetic field, a direct consequence is that all properties of a charged system are periodic functions of the magnetic flux, Φ , the flux periodicity given by the fundamental flux quantum, $\Phi_0 = h/e$. This result is based on the particular combination, $p + eA$, which appears in the Hamiltonian of the system, where p is the momentum, and A the vector potential; here we consider electronic systems, and the charge of an electron is $-e$.

In equilibrium, the system's properties can be calculated from the partition function, which involves a trace over all states of the systems: hence in the classical limit any flux dependence disappears (Bohr-van-Leeuwen-Theorem), and the persistent current, $I(\Phi) = -\partial F(\Phi)/\partial\Phi$, vanishes; $F(\Phi)$ denotes the thermodynamic potential. Thus very small systems and very low temperatures are required for a finite (non-zero) $I(\Phi)$ to exist.

In fact “normal” persistent currents, of the order of a few nA, have been seen in several experiments, for temperatures below 1 K [1–4].

In contrast to the experiments [1–3] which used a SQUID technique in order to detect the magnetic moment induced by the current, the most recent study [4] employed a nano-electromechanical technique: the rings were placed on a cantilever, whose oscillation frequency can be measured with extremely high accuracy. The perimeter of the studied rings varied between 0.6 and 1.6 μm .

Assuming time reversal invariance, the Fourier expansion of the persistent current is given by

$$I(\Phi) = I_1 \sin(2\pi\Phi/\Phi_0) + I_2 \sin(4\pi\Phi/\Phi_0) + \dots$$

Because of the disorder always present in small rings, due to the fabrication process, the amplitudes I_1 and I_2 are random quantities. For example, I_1 fluctuates from ring to ring; in particular, it changes its sign, hence the average is expected (and was found) to be zero. The size of $\langle(I_1)^2\rangle^{1/2}$, on the other hand, has been debated for many years; the theoretical prediction $\langle(I_1)^2\rangle^{1/2} \approx E_c/\Phi_0$ [5] was now convincingly confirmed [4]; here $E_c = \hbar D/L^2$ is called Thouless energy, D is the diffusion constant, and L the perimeter of the ring. Concerning the second harmonic, I_2 , it was pointed out rather early [6] that the effective electron-electron interaction gives an important contribution, which nevertheless is too small compared to the experimental result [1]. A recent paper discusses the question whether a small amount of paramagnetic impurities can resolve this discrepancy [7]. For an introduction into persistent currents, see [8].

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